



THE UNIVERSITY OF
SYDNEY



Can we tame earthquakes?

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David Berry / Getty Images



Anti-lock Braking System



Earthquakes and faults



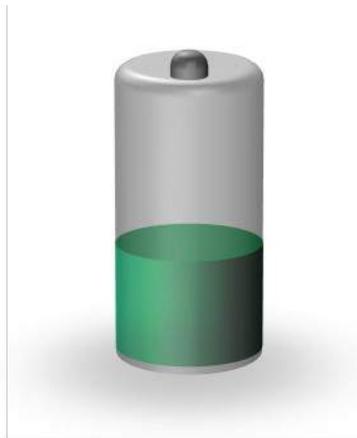
Greendale fault, New Zealand



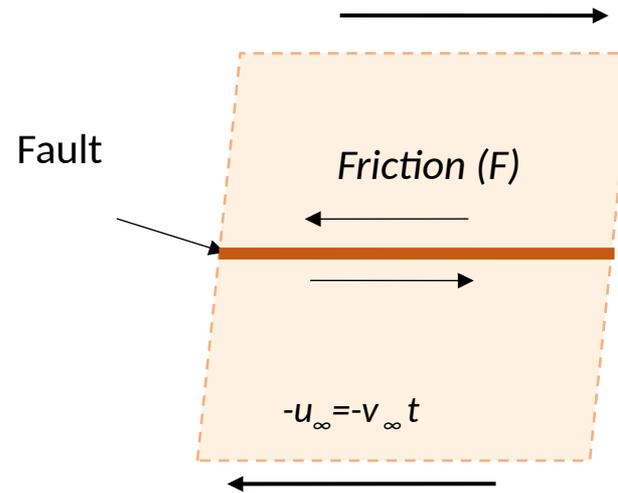
Greendale fault, New Zealand

Earthquakes in simple words

Earthquakes in simple words



Far field tectonic displacements ($\sim\text{cm/yr}$), $u_\infty = v_\infty t$

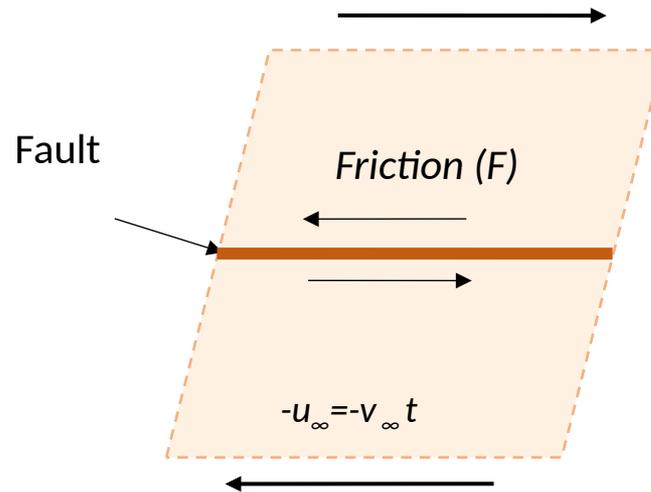


Storage of elastic energy in the crust $U_{el} \uparrow$

Earthquakes in simple words



Far field tectonic displacements ($\sim\text{cm/yr}$), $u_\infty = v_\infty t$

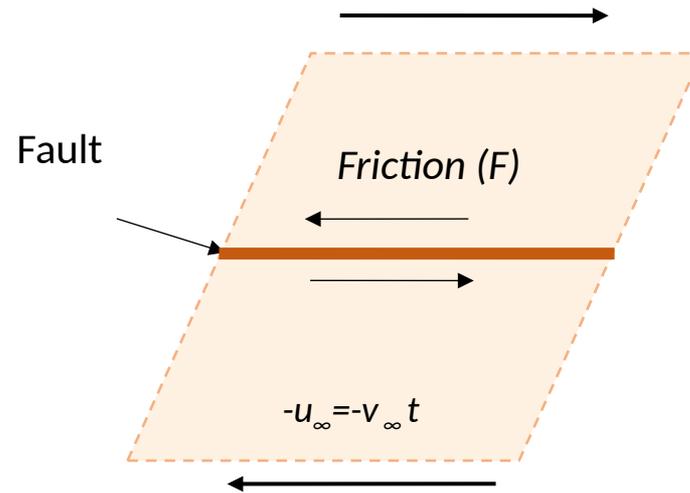


Storage of elastic energy in the crust $U_{el} \uparrow$

Earthquakes in simple words



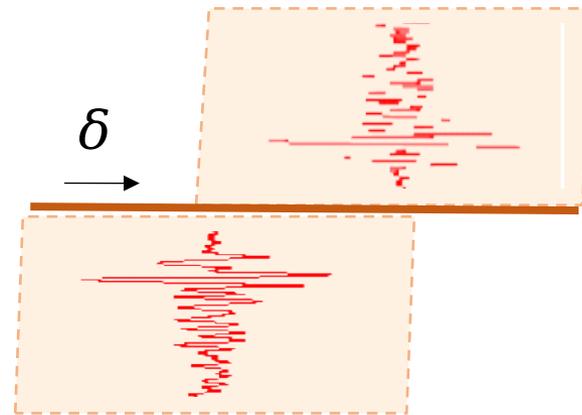
Far field tectonic displacements ($\sim\text{cm/yr}$), $u_\infty = v_\infty t$



Storage of elastic energy in the
crust $U_{el} \uparrow$

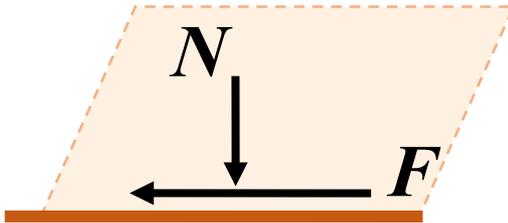
Earthquakes in simple words

Instability



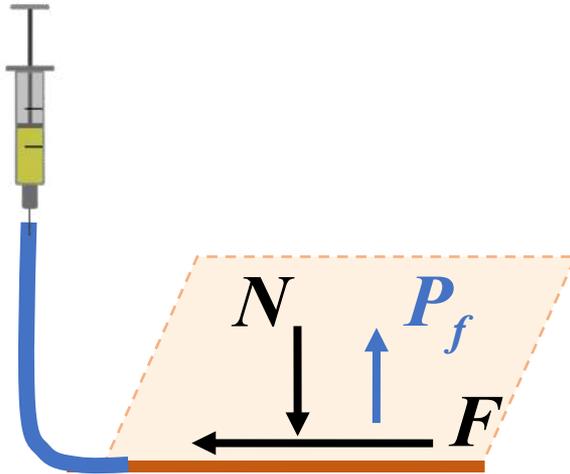
Injecting fluids and friction

$$F = \mu \cdot N$$



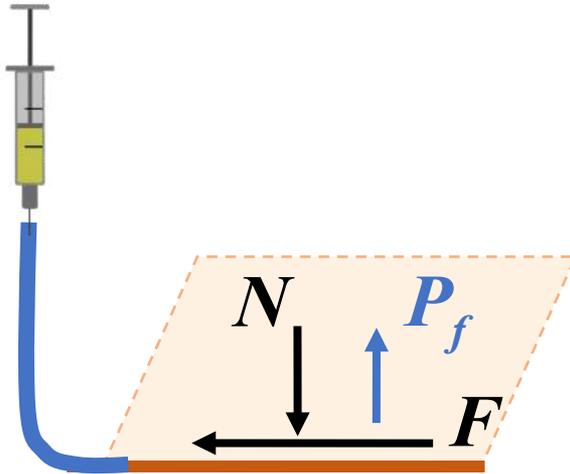
Injecting fluids and friction

$$F = \mu \cdot (N - P_f)$$



Injecting fluids and friction

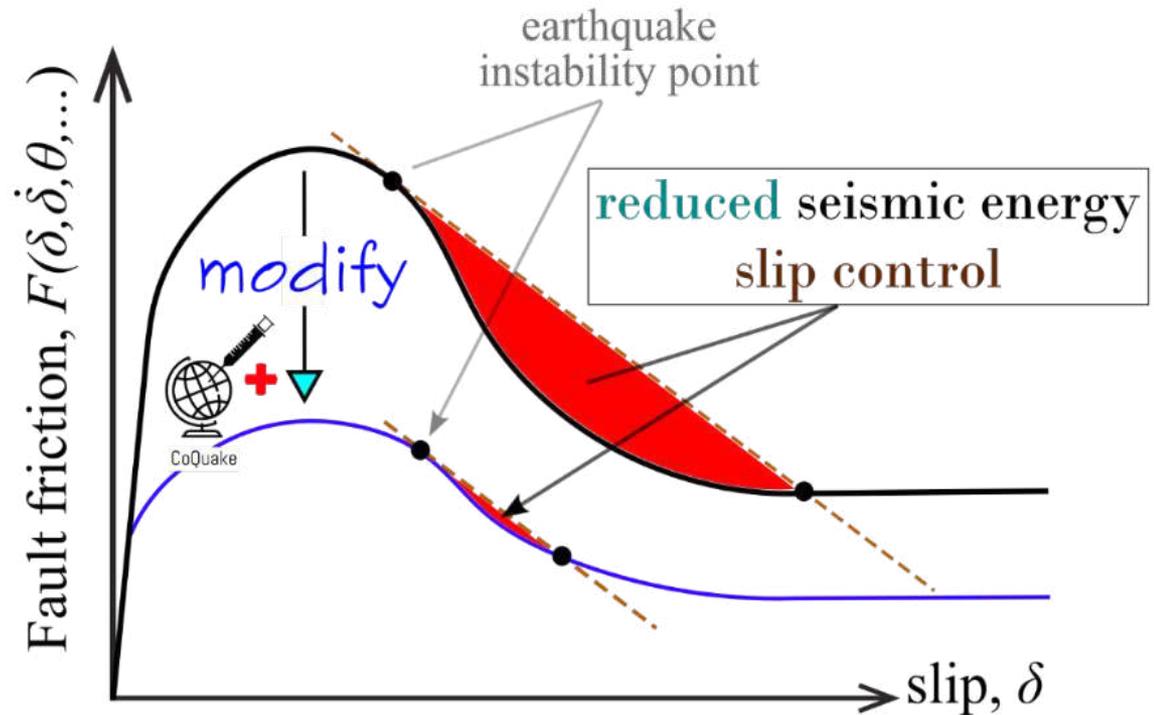
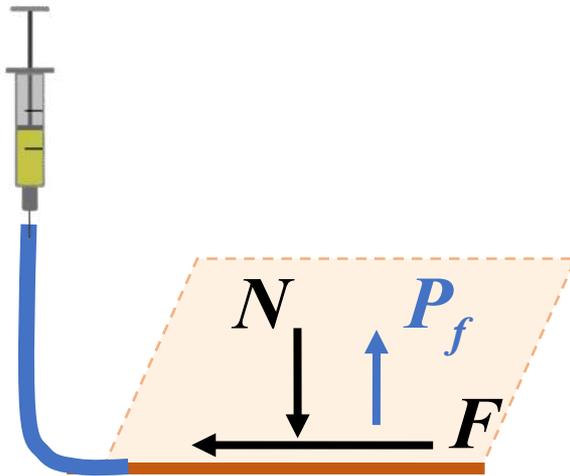
$$F = \mu \cdot (N - P_f)$$



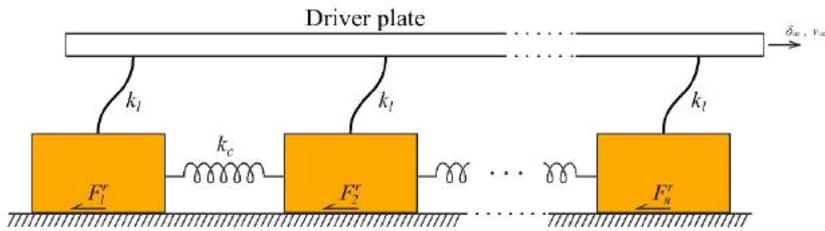
Injecting fluids and friction

$$F = \mu \cdot (N - P_f)$$

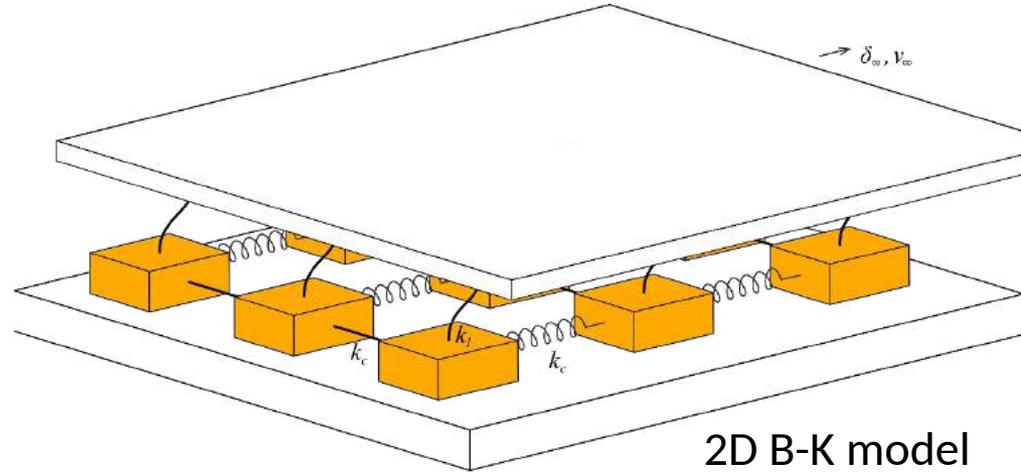
[Terzaghi, 1925]



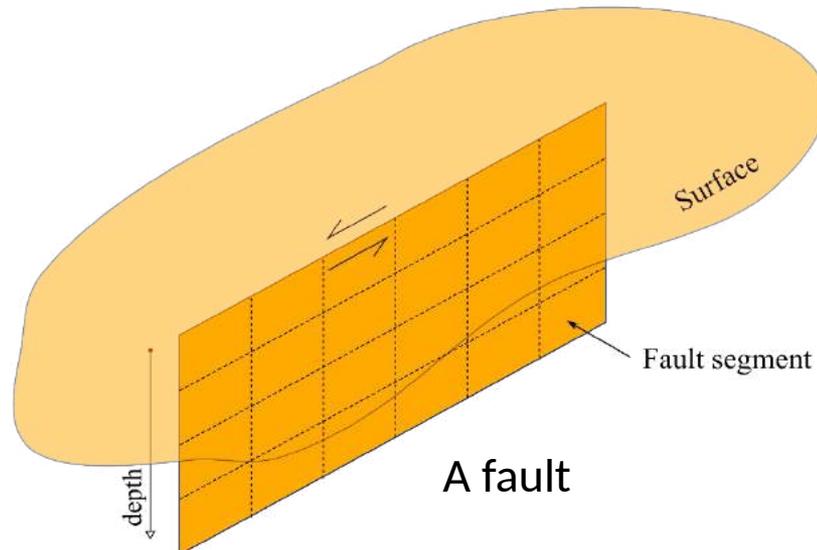
Increasing complexity: Burridge-Knopoff & Faults



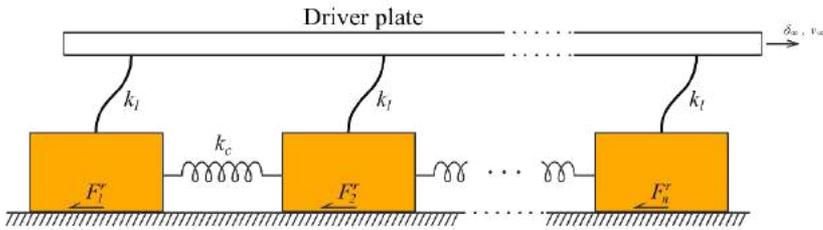
1D B-K model



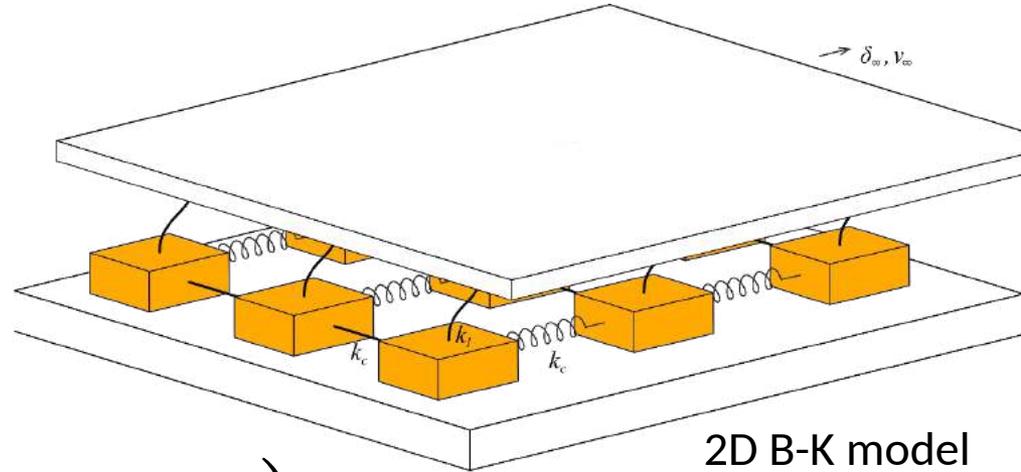
2D B-K model



Increasing complexity: Burridge-Knopoff & Faults

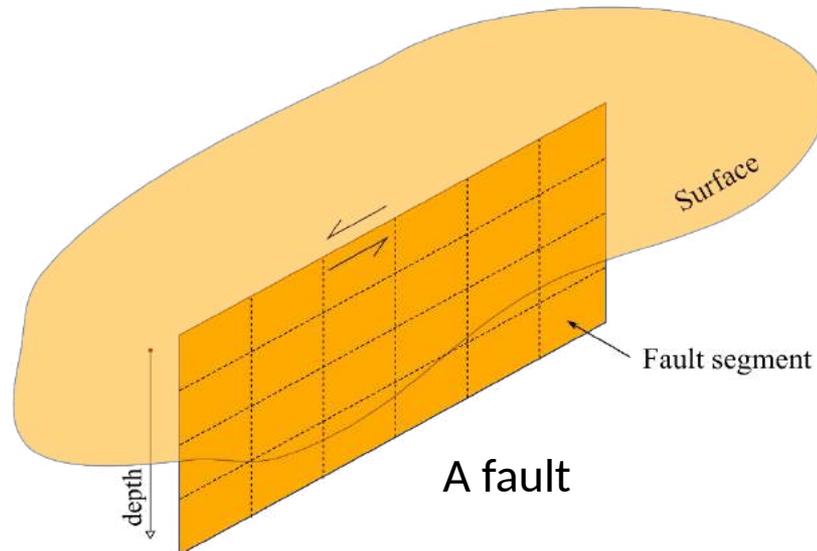


1D B-K model

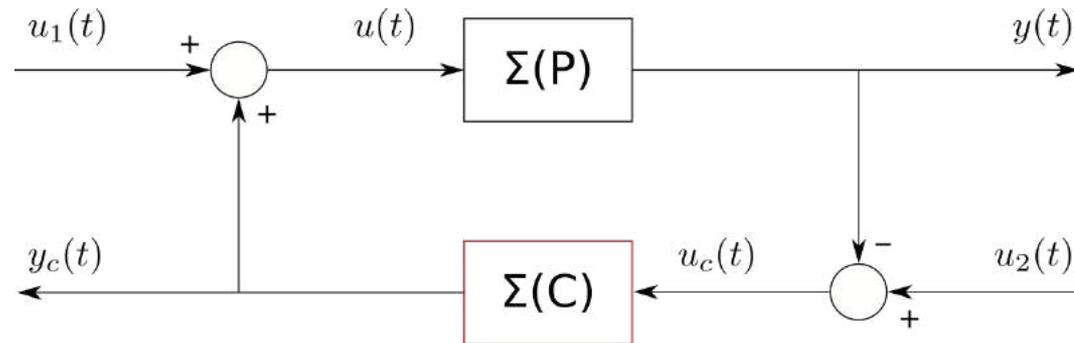


2D B-K model

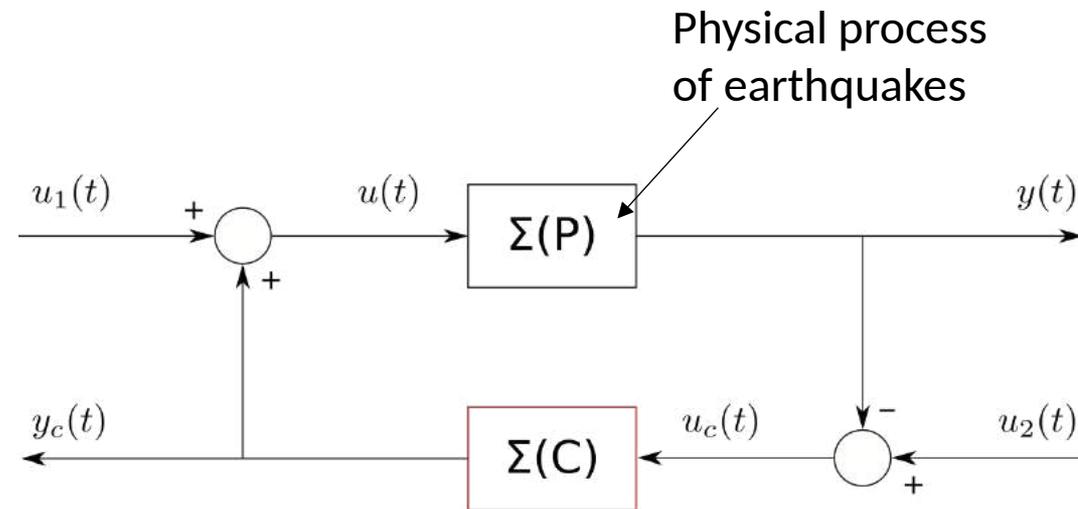
$$\dot{x} = f(x, p_f, t)$$



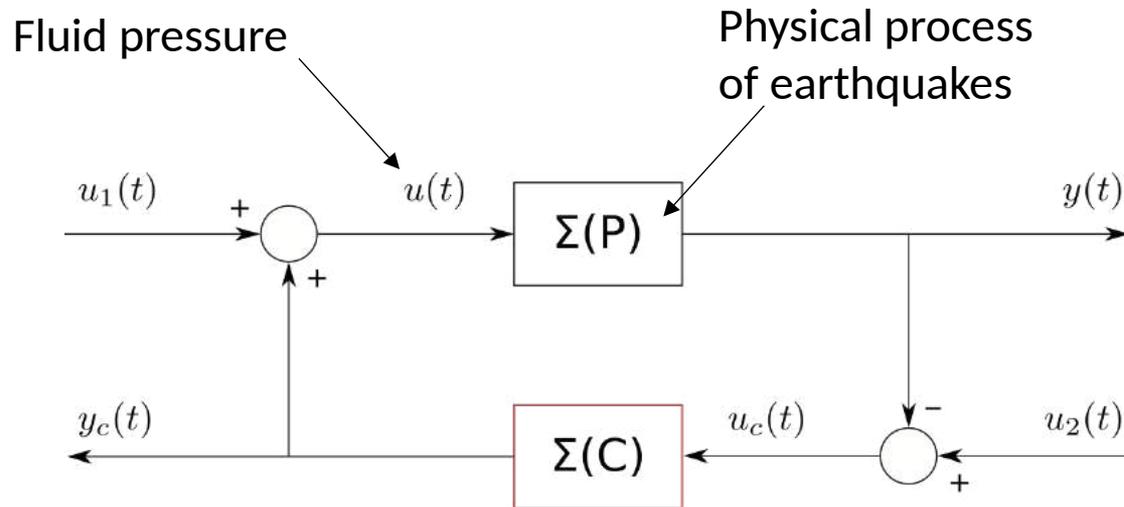
The mathematical theory of control



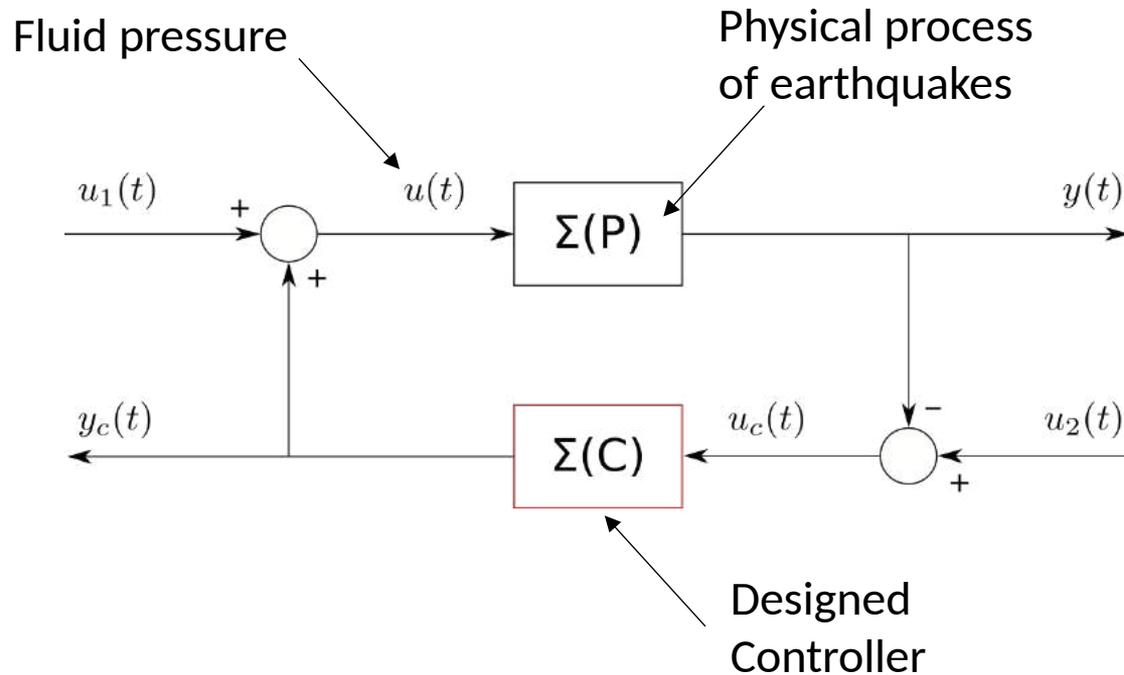
The mathematical theory of control



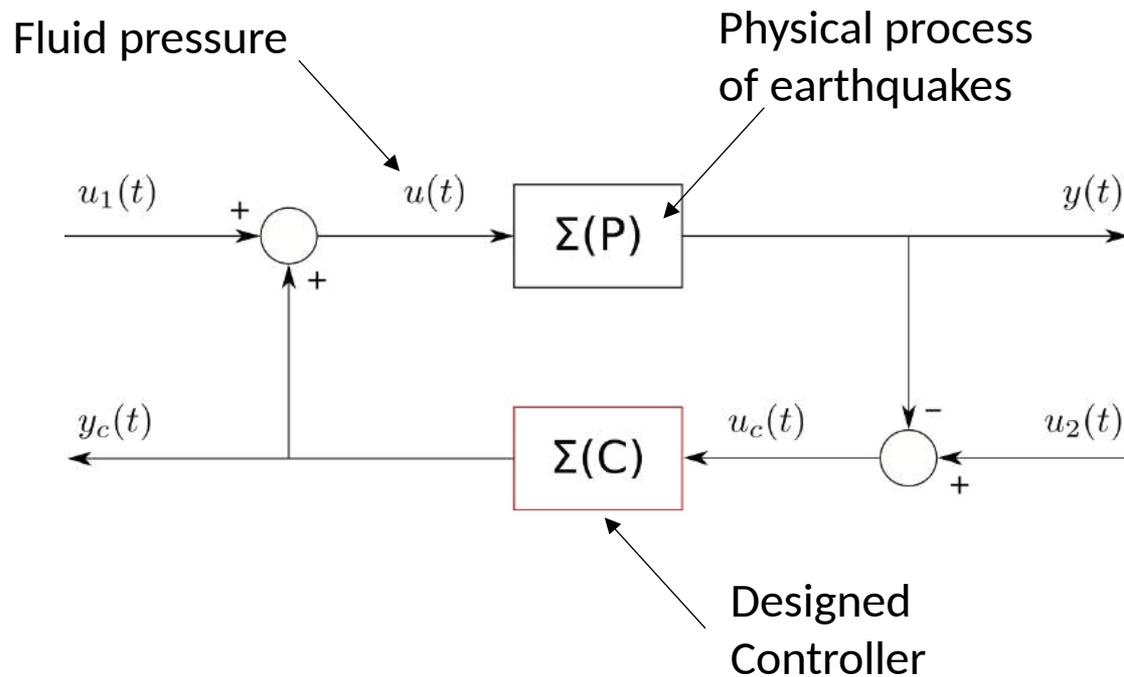
The mathematical theory of control



The mathematical theory of control



The mathematical theory of control



Target: stabilization, tracking,
optimization & robustness

Robust non-linear control

If:

1) The friction coefficient is Lipschitz continuous w.r.t the states x :

$$|\mu(x, t) - \mu(0)| < \beta |x|, \beta > 0$$

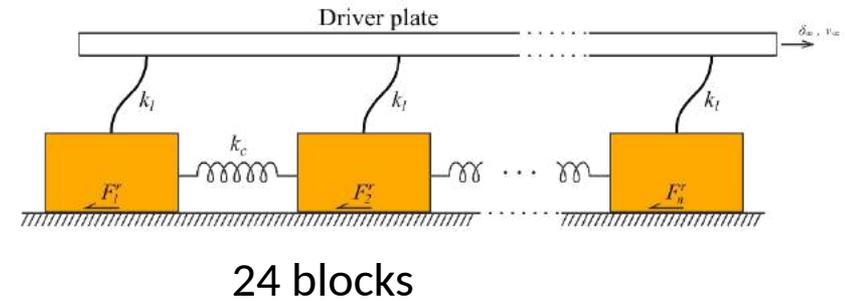
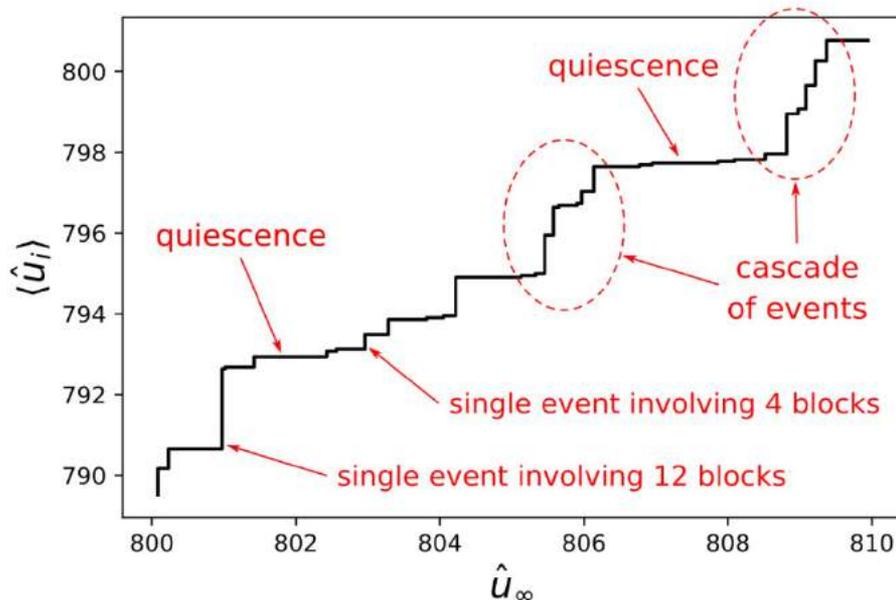
2) The friction coefficient has a lower bound: $\mu(x, t) > c > 0$

3) Elasticity and viscosity of the surrounding rocks are bounded

4) Diffusivity has a lower bound greater than zero

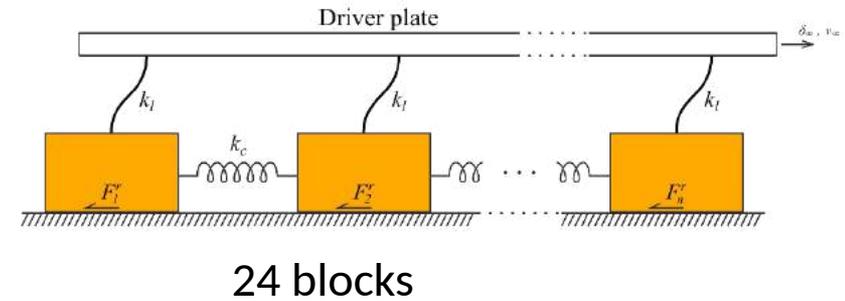
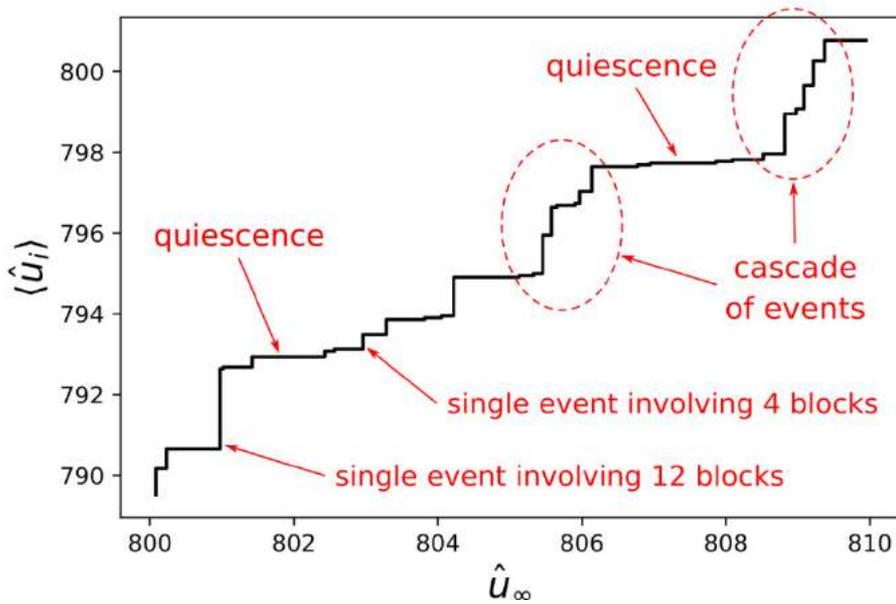
then we can design an output feedback stabilizing controller and we can achieve asymptotic tracking.

Example of control of a complex frictional system

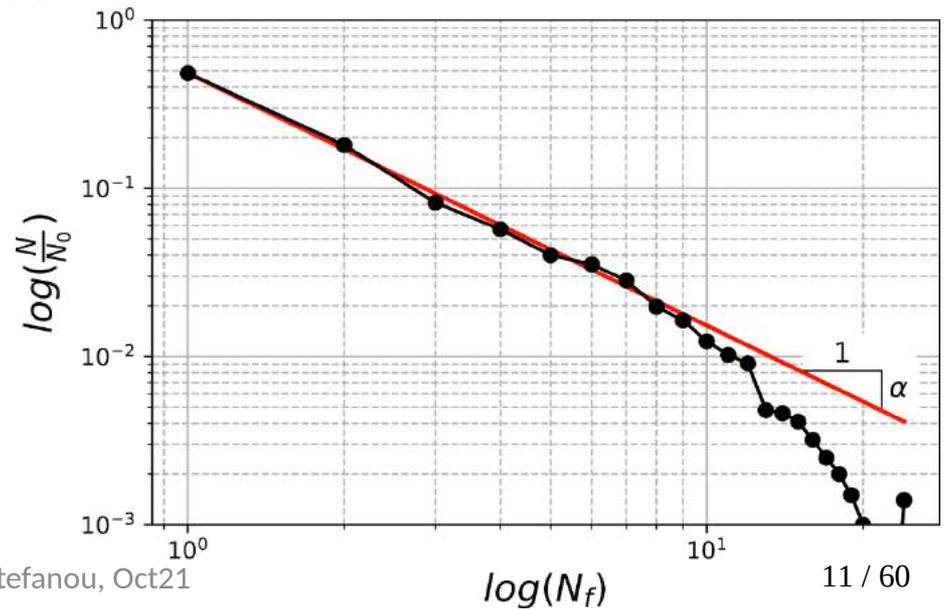


[see also Turcotte, 1992,
– among many others]

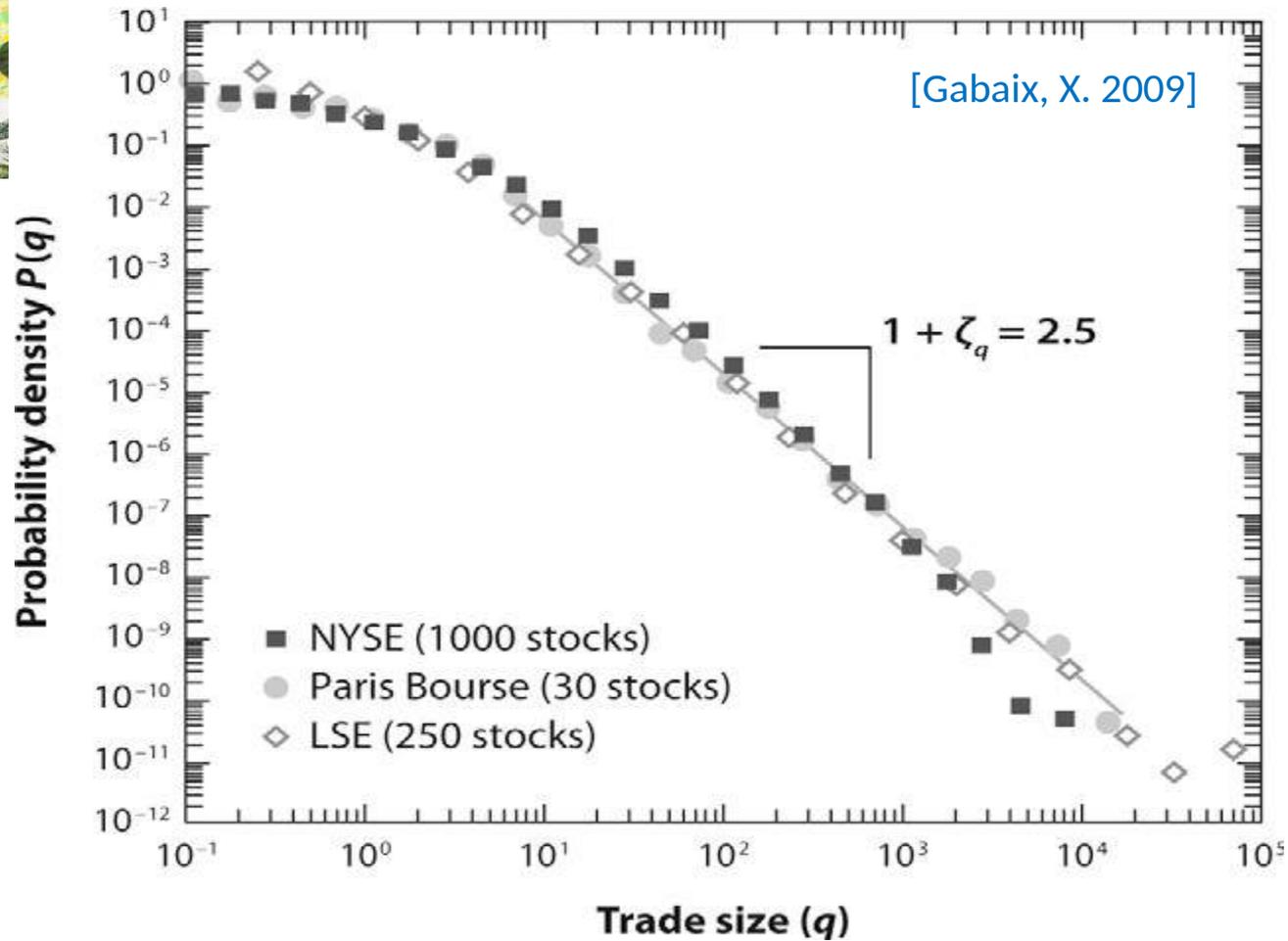
Example of control of a complex frictional system



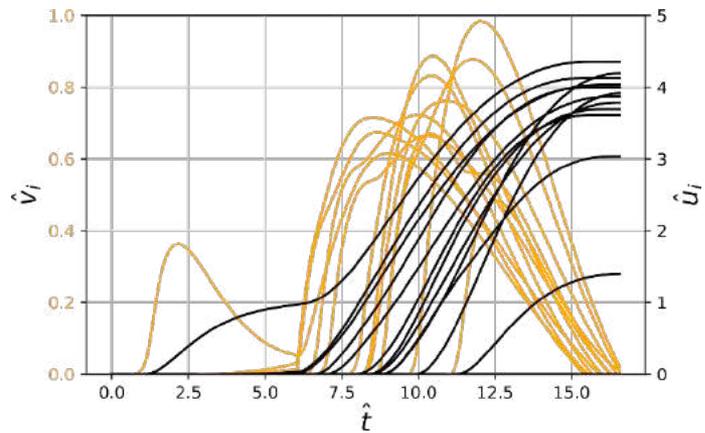
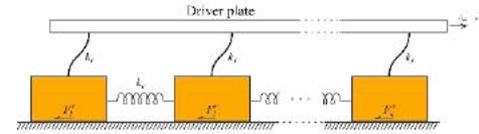
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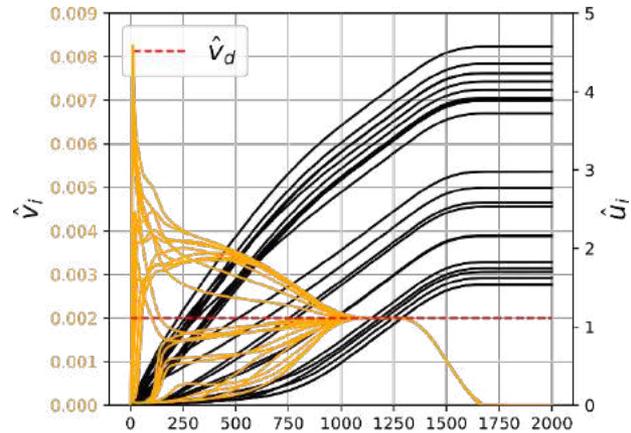
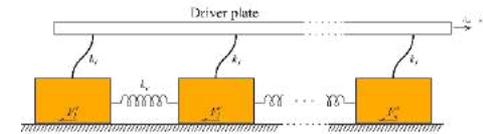
Stock markets?

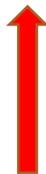


Tracking and global stabilization

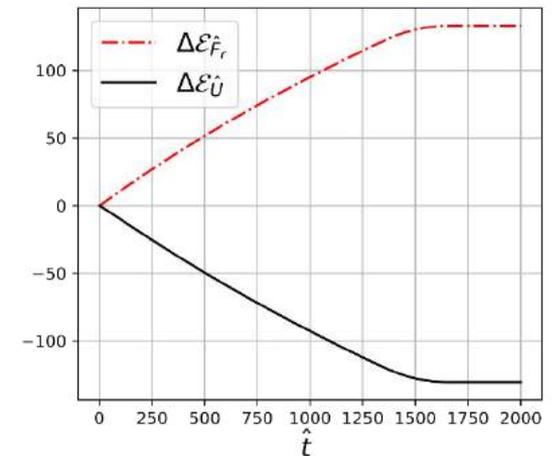
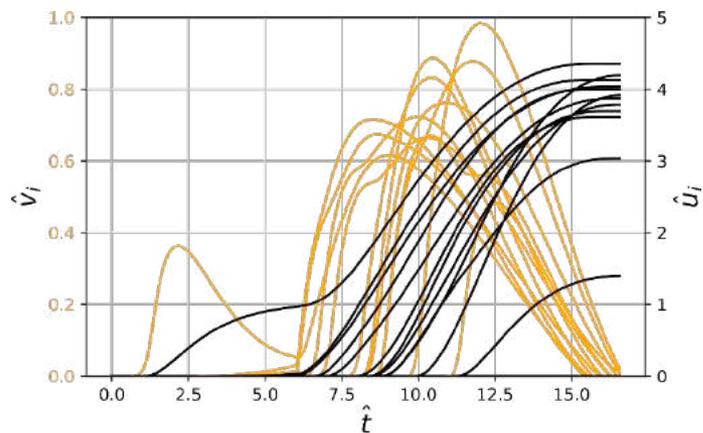
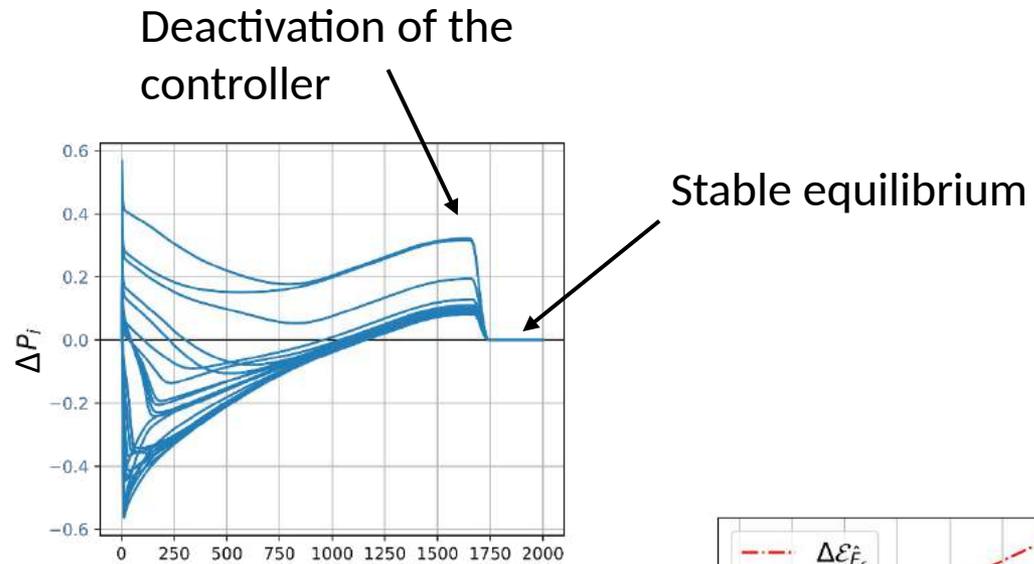


Tracking and global stabilization





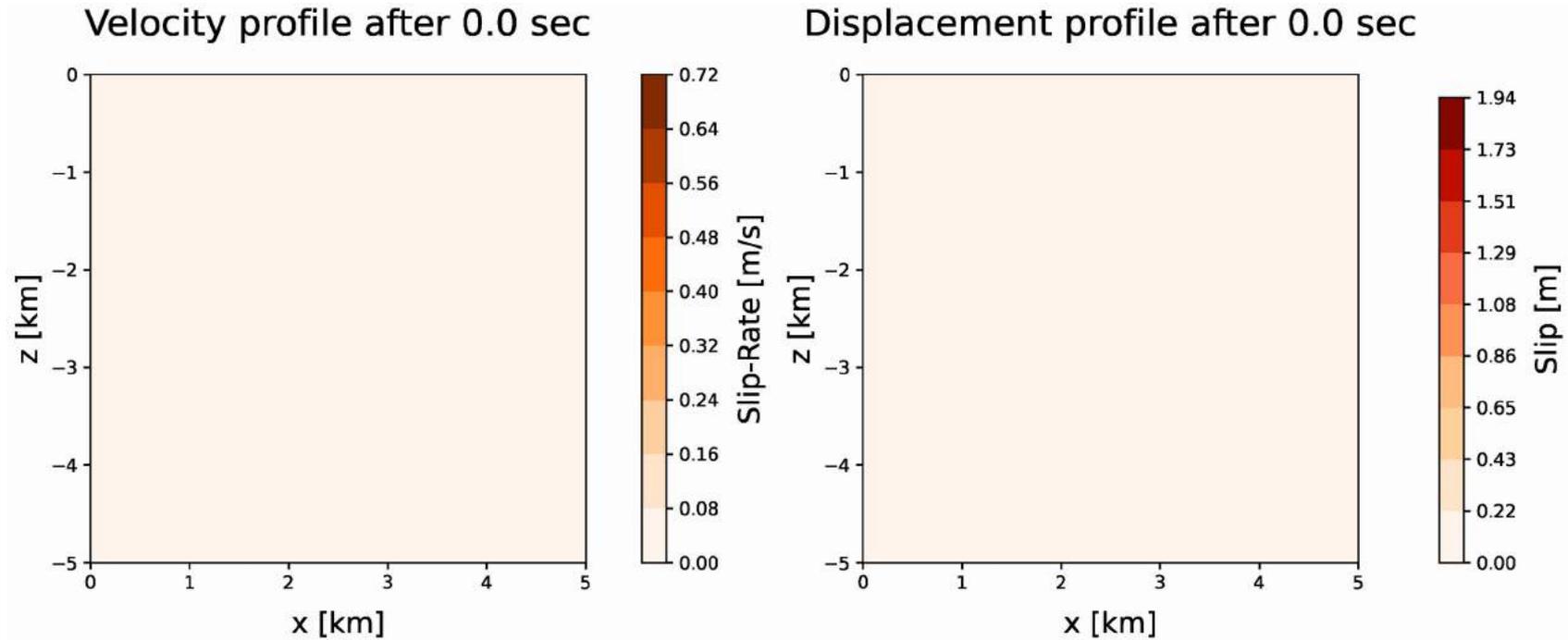
Control



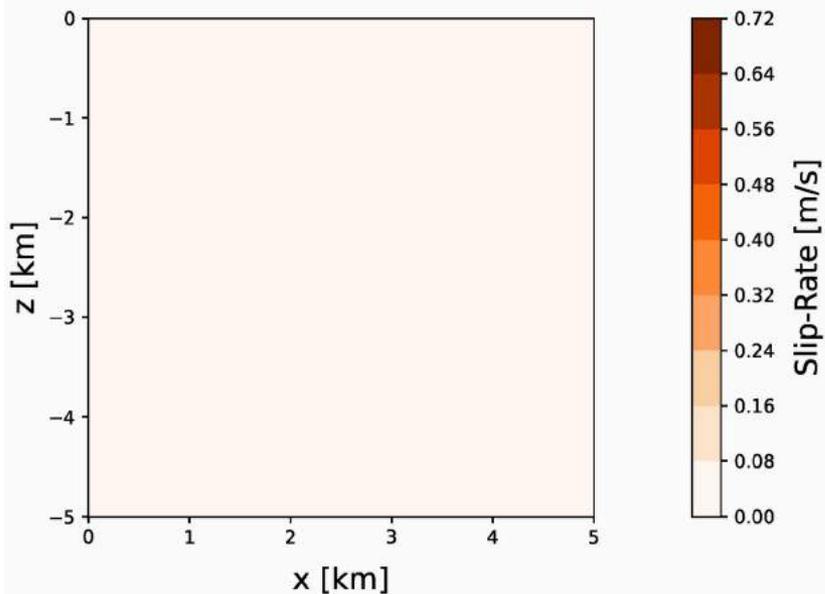
Strike-slip fault with Rate & State friction

a, b, d_c distributed over the fault area according to a log-normal distribution

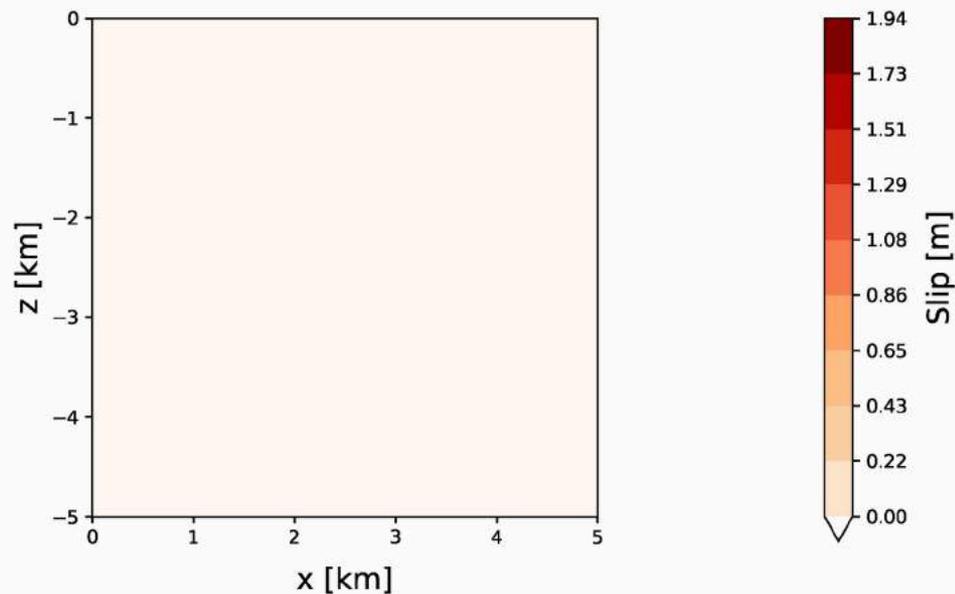
Without control (open loop)



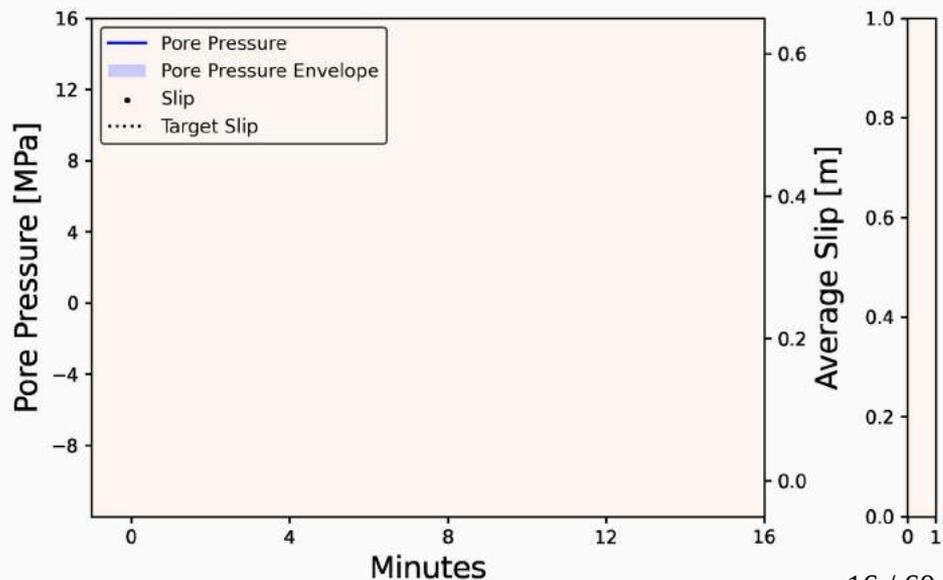
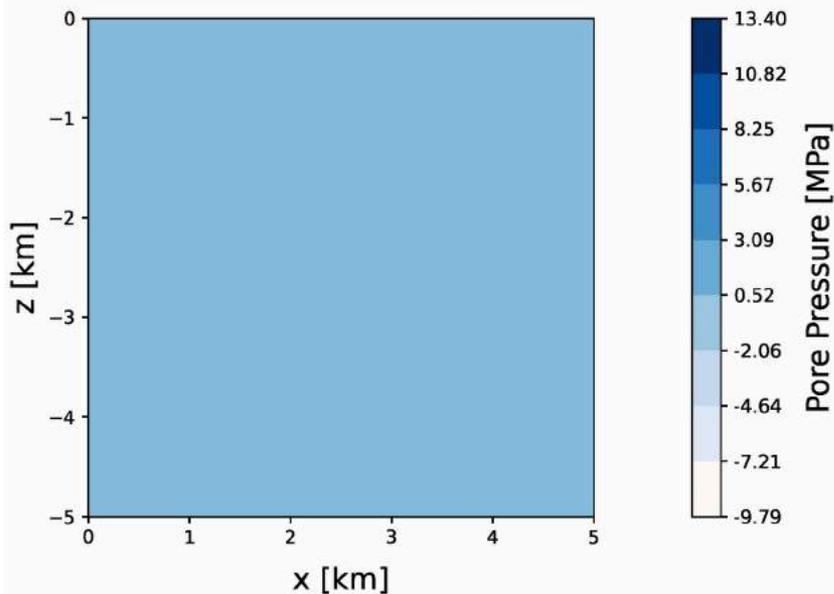
Velocity profile after 0.0 min



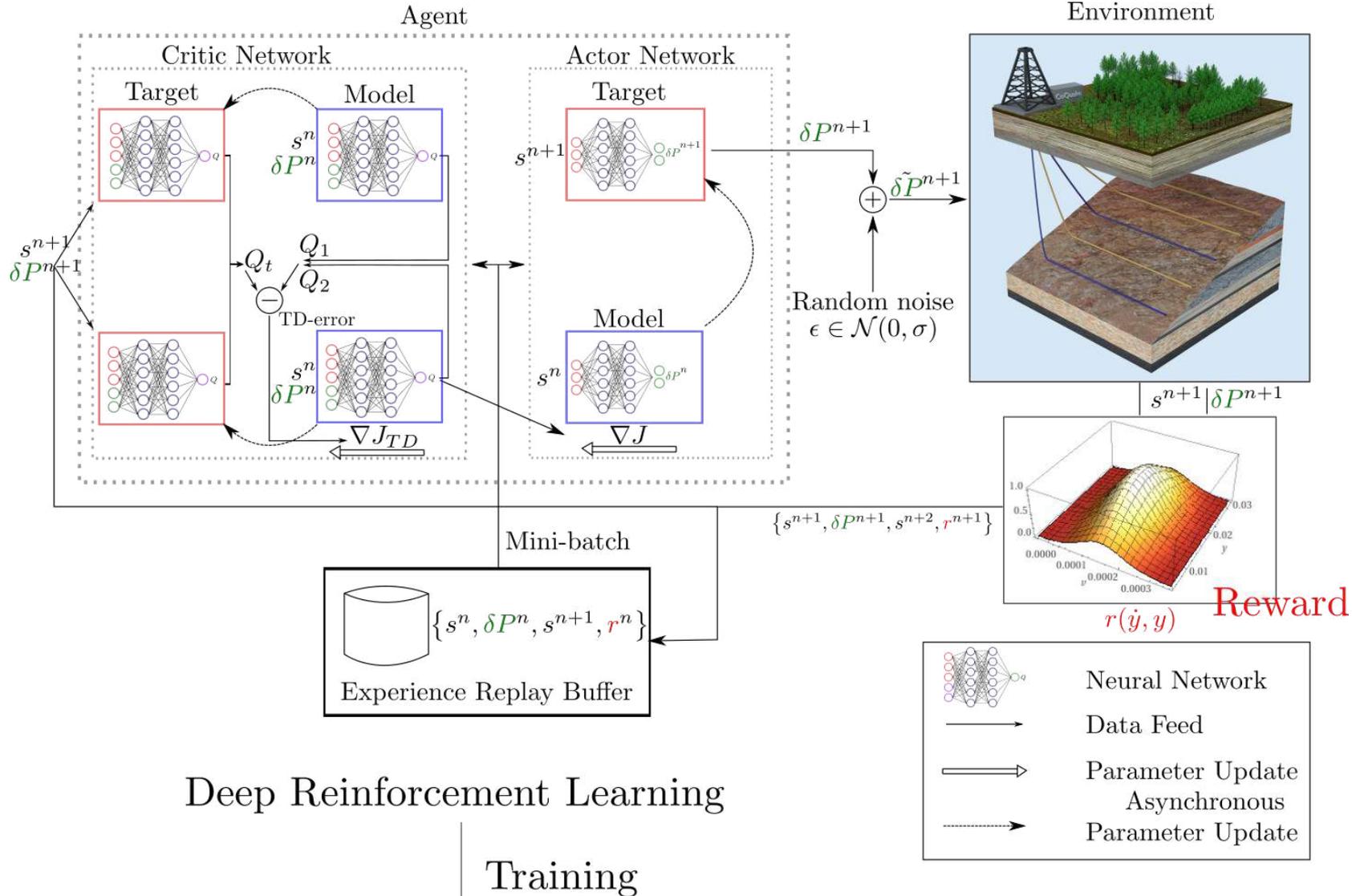
Displacement profile after 0.0 min



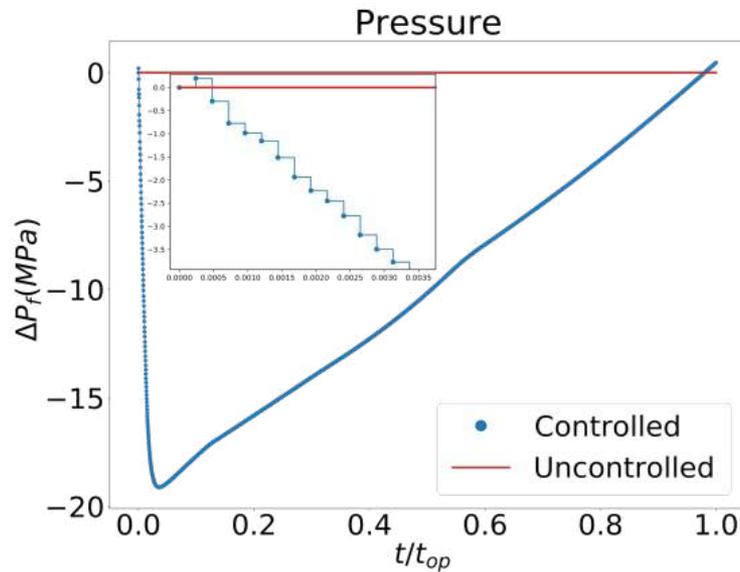
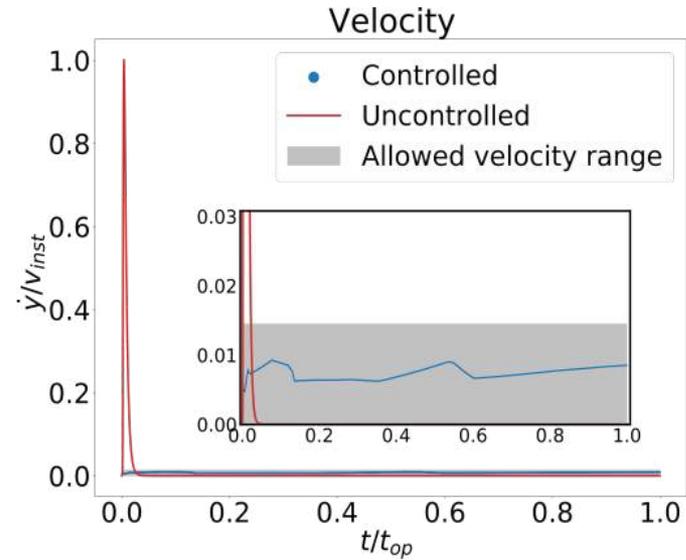
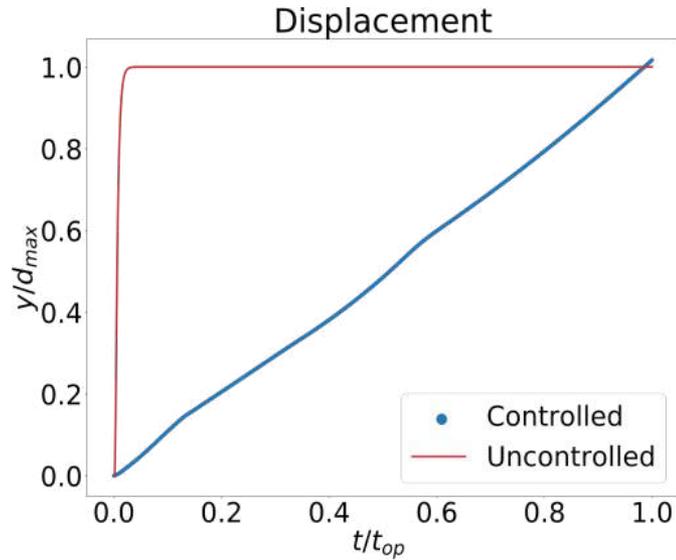
Pore-pressure profile after 0.0 min



Earthquake control with Reinforcement Learning

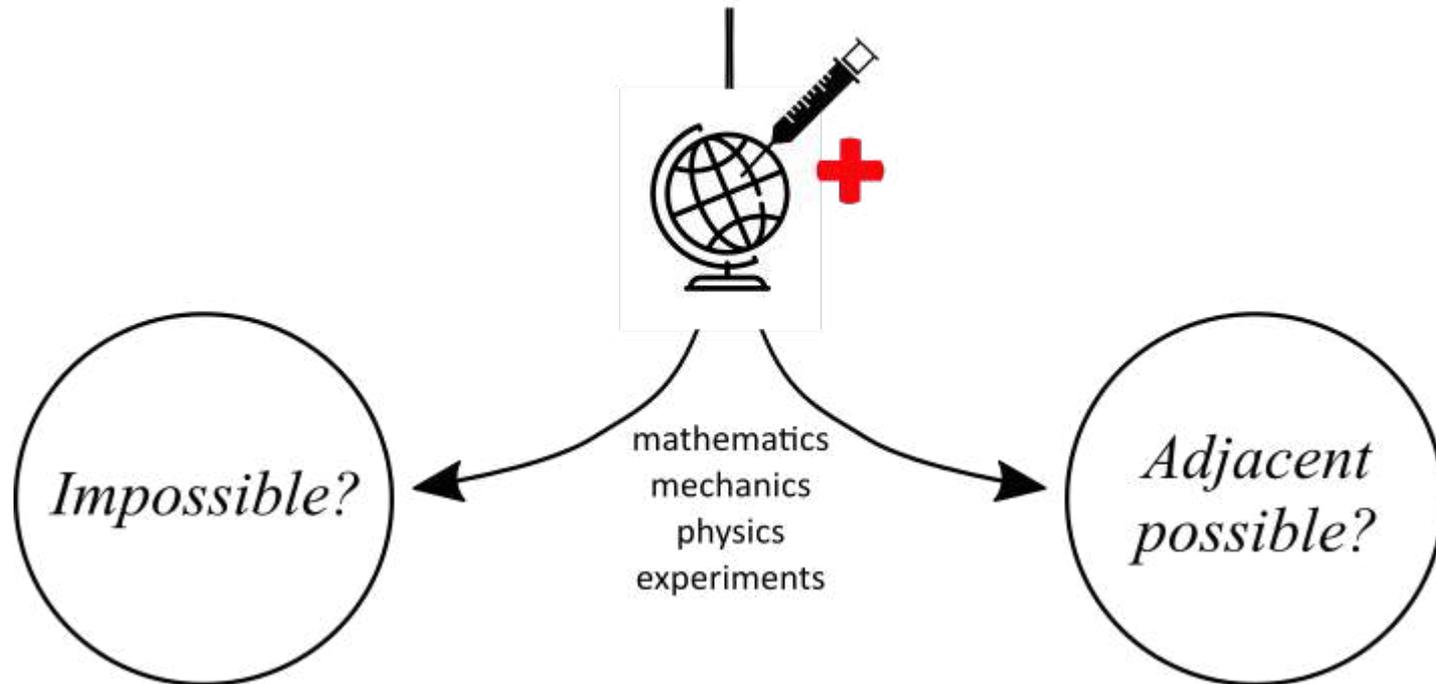


DRL, discrete time dynamics & results



Conclusions

Earthquake Control?



Thank you!



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CoQuake.com



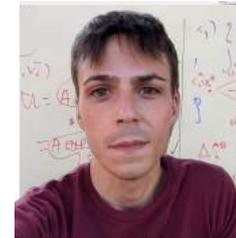
George
Tzortzopoulos-
Marinis



Philipp Braun



Alexandros
Stathas



Filippo Masi



Timos
Papachristos

ioannis.stefanou@ec-nantes.fr

TED^x Talk: <https://youtu.be/9JXdv-3e2bc>

www.coquake.eu

www.blastructures.eu



Diego
Gutiérrez-
Oribio



Pravin
Badarayani

Thank you!



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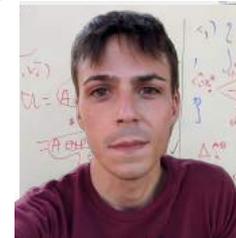
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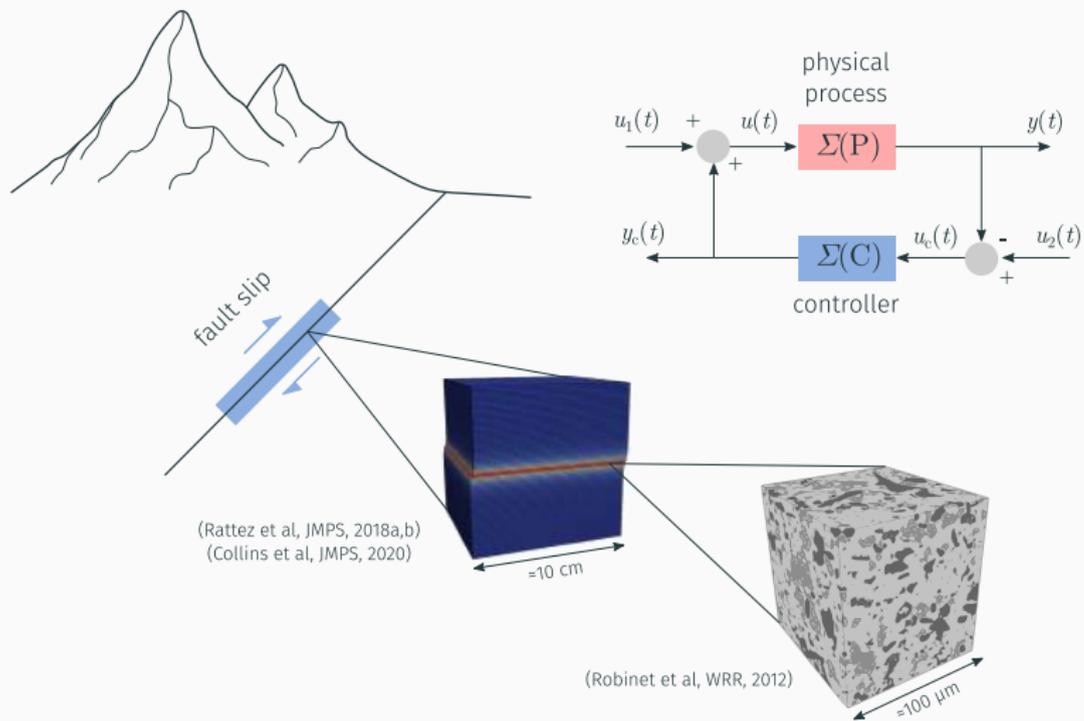
Ecole Centrale de Nantes, GeM (Institut de Recherche en Génie Civil et Mécanique), France

ioannis.stefanou@ec-nantes.fr filippo.masi@ec-nantes.fr

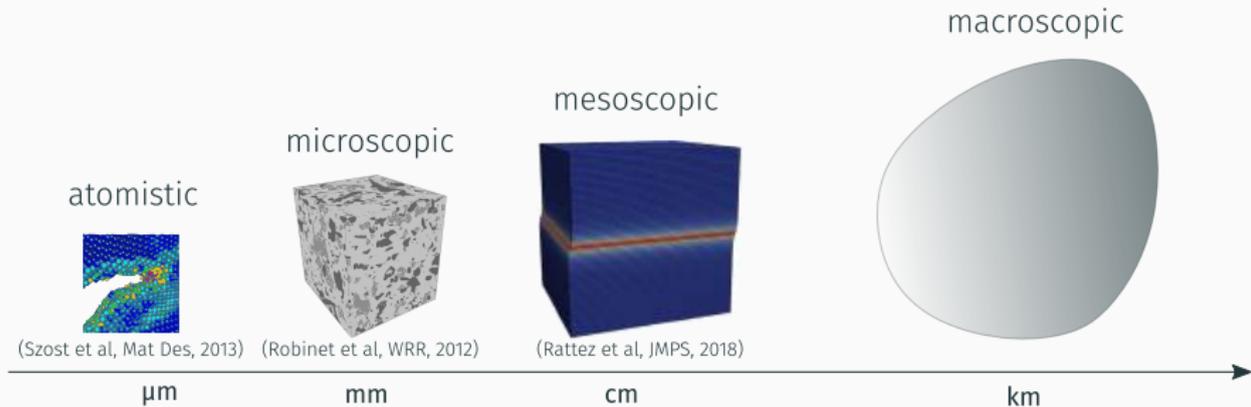


BRIDGING THE SCALES

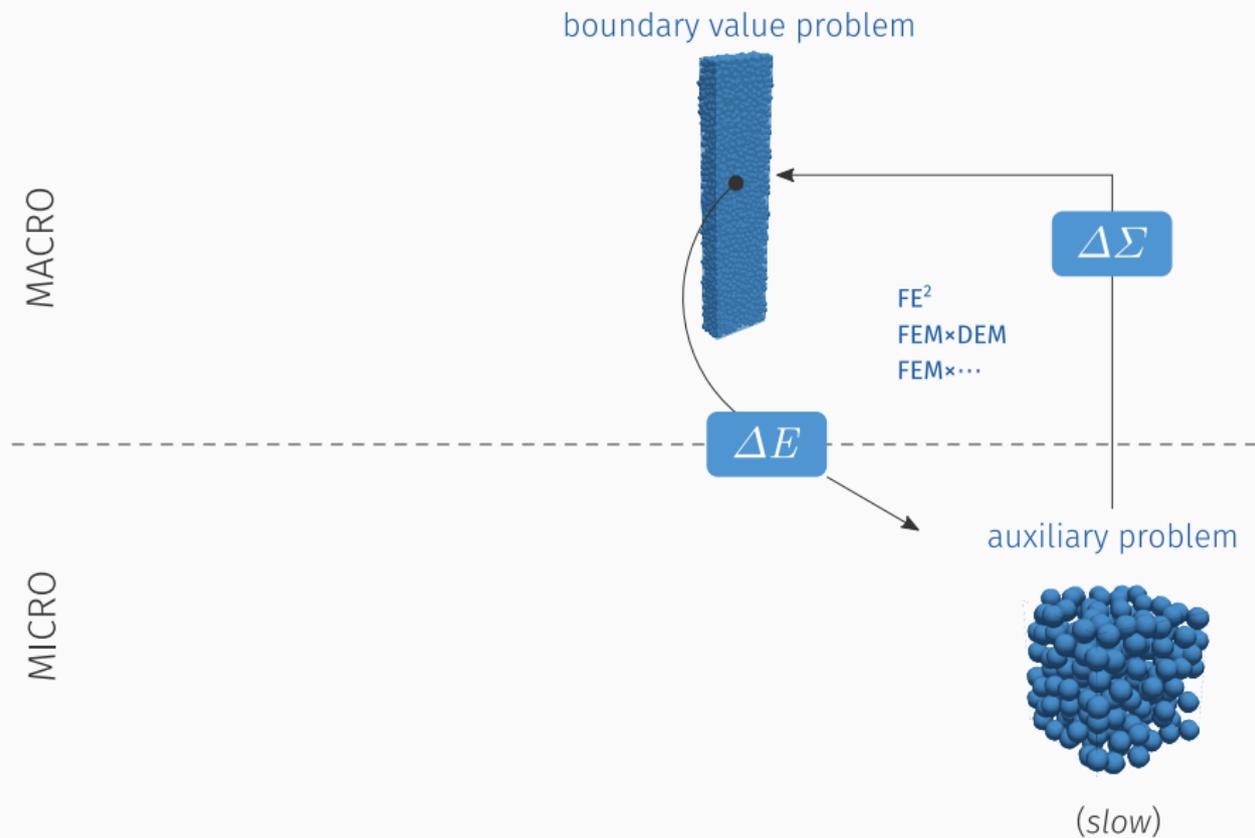
Systems involving complex materials – multiple inherent scales



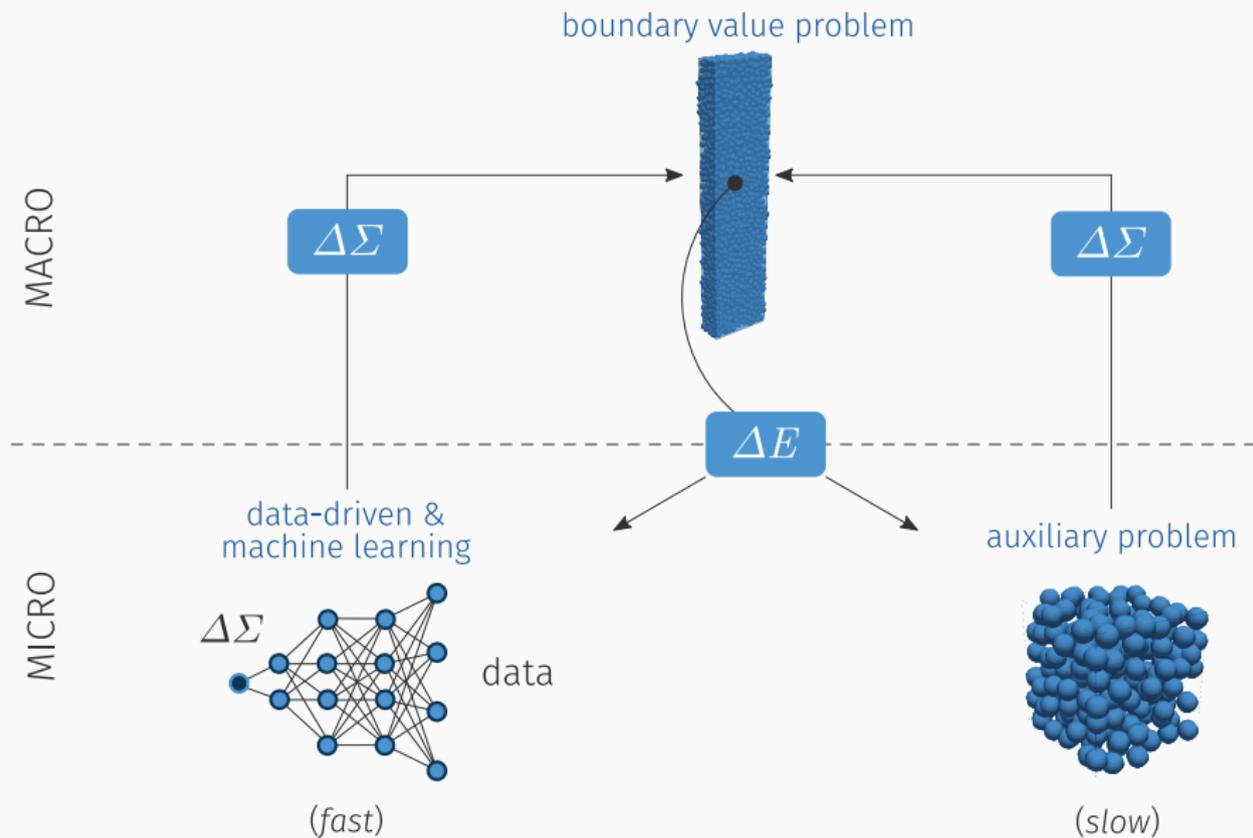
BRIDGING THE SCALES – MULTISCALE MODELING



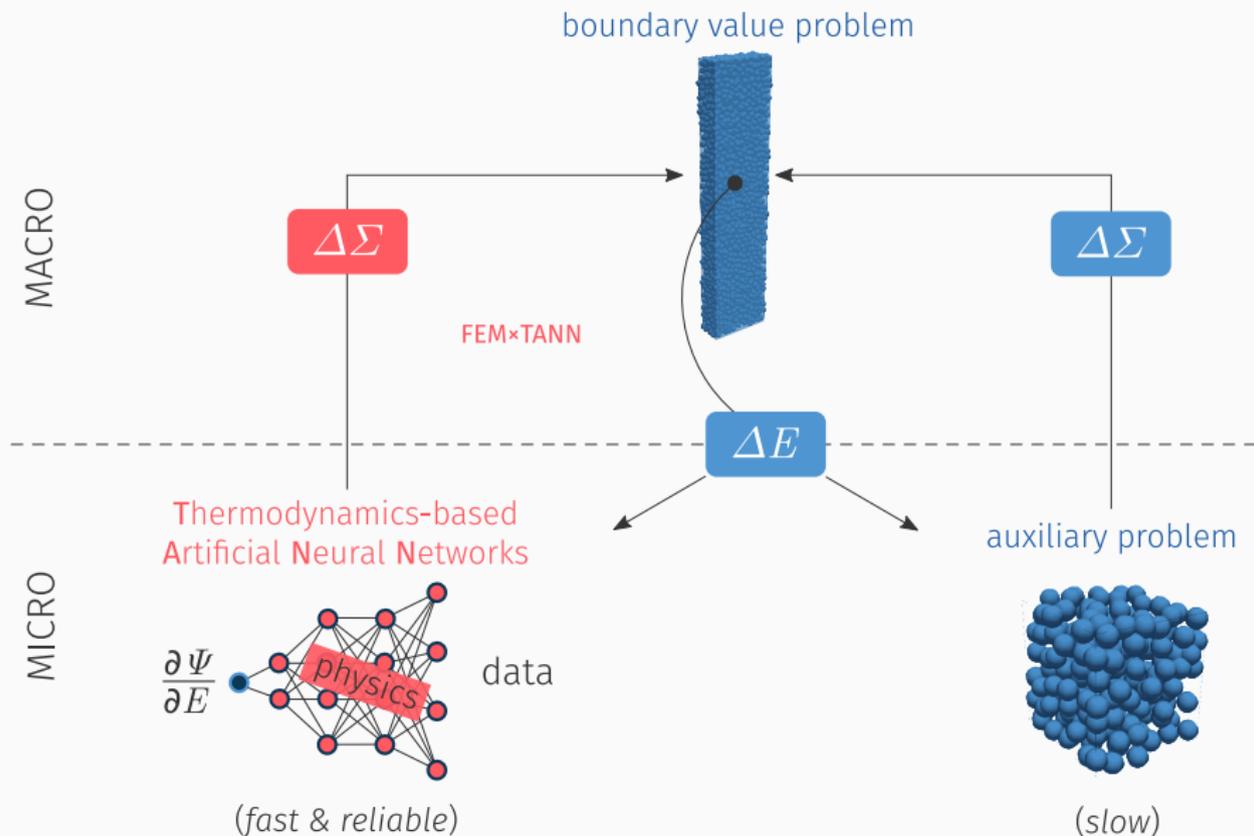
BRIDGING THE SCALES – MULTISCALE MODELING



BRIDGING THE SCALES – DATA-DRIVEN MODELING



BRIDGING THE SCALES – TANN



THERMODYNAMICS

Clausius-Duhem inequality (local)

$$d = s : \dot{f} - \dot{\psi} - \eta \dot{\theta} \geq 0$$

d mechanical dissipation rate

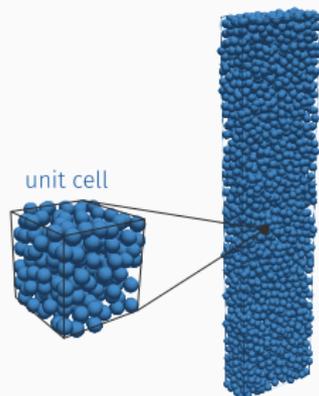
s 1st Piola-Kirchhoff stress

f deformation gradient

ψ free-energy

η entropy

θ absolute temperature



Clausius-Duhem inequality (volume average)

$$D = S : \dot{F} - \dot{\Psi} - H^\dagger \dot{\Theta} \geq 0$$

with $Y = \langle y \rangle = \frac{1}{|V|} \int_V y dx$ and $\langle \eta \dot{\theta} \rangle = H^\dagger \dot{\Theta}$

THERMODYNAMICS

$$\text{STATE SPACE}(t) = \{\chi(t), Z(t)\}$$

state variables: **observable** $\chi = \{\Theta, F\}$ and **internal** $Z = \{?\}$

state functions: $\Psi(\chi, Z)$ $S(\chi, Z)$ $H^\dagger(\chi, Z)$ $D(\chi, Z, \dot{Z})$

time differentiation: $\dot{\Psi} = \frac{\partial \Psi}{\partial \Theta} \dot{\Theta} + \frac{\partial \Psi}{\partial F} : \dot{F} + \sum_{k=1}^{N_{ISV}} \frac{\partial \Psi}{\partial Z_k} \cdot \dot{Z}_k$

THERMODYNAMIC RESTRICTIONS:

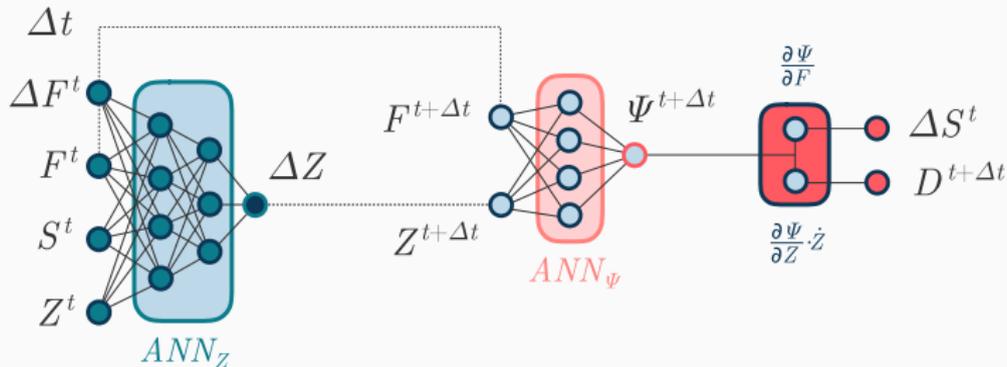
$$\text{stress: } S = \frac{\partial \Psi}{\partial F}$$

$$\text{dissipation rate: } D = - \sum_{k=1}^{N_{ISV}} \frac{\partial \Psi}{\partial Z_k} \cdot \dot{Z}_k \geq 0$$

$$\text{entropy: } H^\dagger = - \frac{\partial \Psi}{\partial \Theta}$$

THERMODYNAMICS-BASED ARTIFICIAL NEURAL NETWORKS

$$\text{state at } t + \Delta t = \text{TANN}\{\text{state at } t, \Delta F^t\}$$



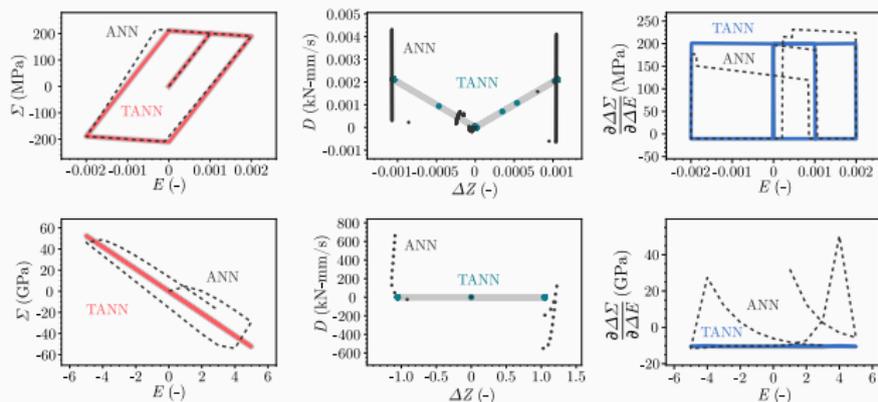
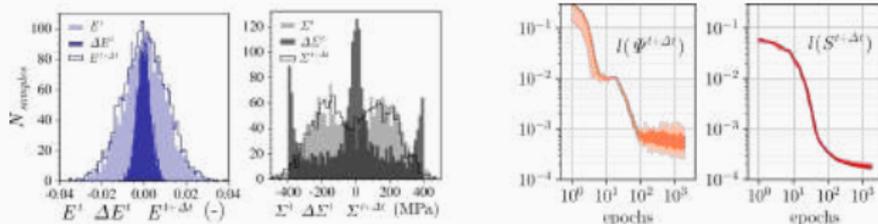
$$\mathcal{L}(o) = \frac{1}{N} \sum_{i=1}^N \ell(o_i)$$

$$\ell(o_i) = \|\bar{o}_i - o_i\|_1$$

$$\ell = \lambda^{\dot{Z}} \ell(\dot{Z}) + \lambda^{\Psi} \ell(\Psi) + \lambda^{\Delta S} \ell(\Delta S) + \lambda^D \ell(D)$$

(Masi et al. J Mech Phys Solids 147, 2021; Masi et al. SPIGL, Springer, 2020)

HYPER- AND HYPO-PLASTICITY

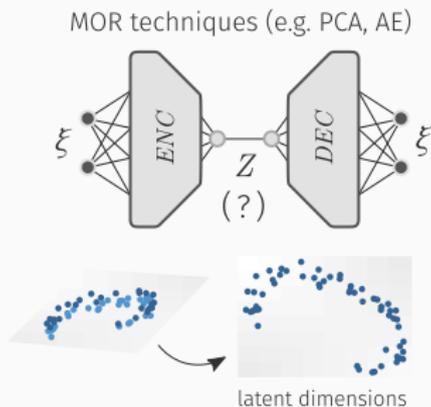
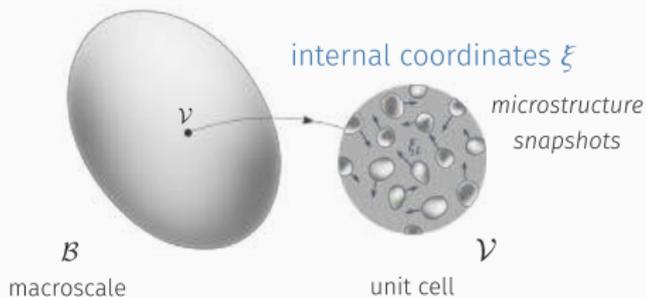


(Masi et al. J Mech Phys Solids 147, 2021; Masi et al. SPIGL, Springer, 2020)

STATE VARIABLES?

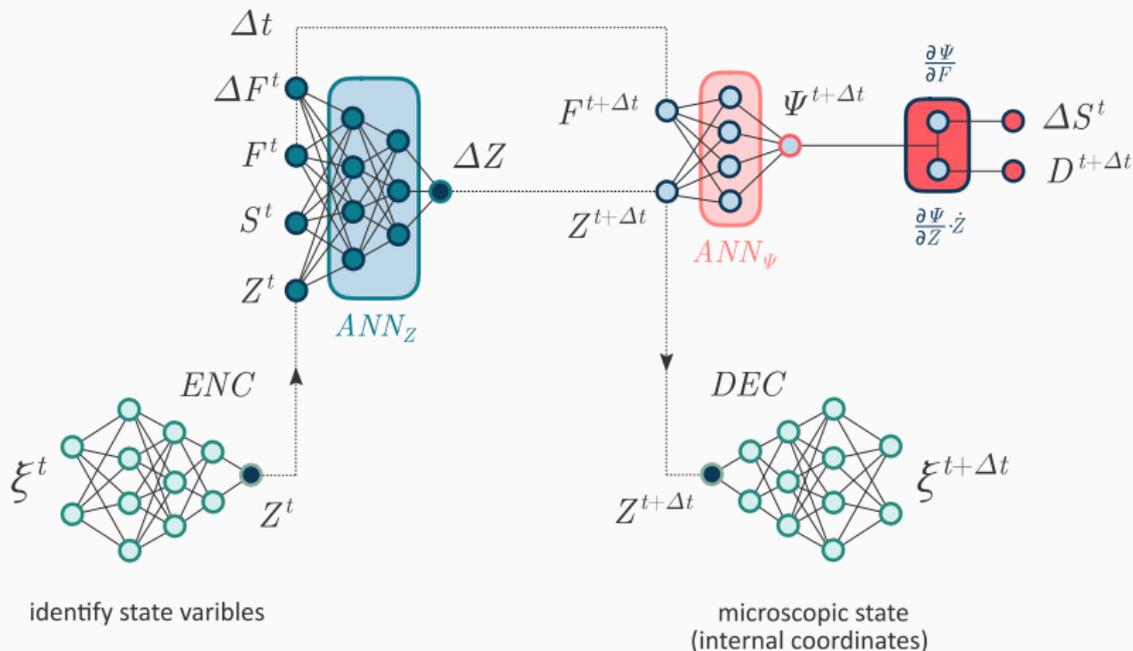
internal coordinates

microscopic state space: displacement, velocity, momentum fields, etc.



(Masi et al. arXiv preprint arXiv:2108.13137, 2021)

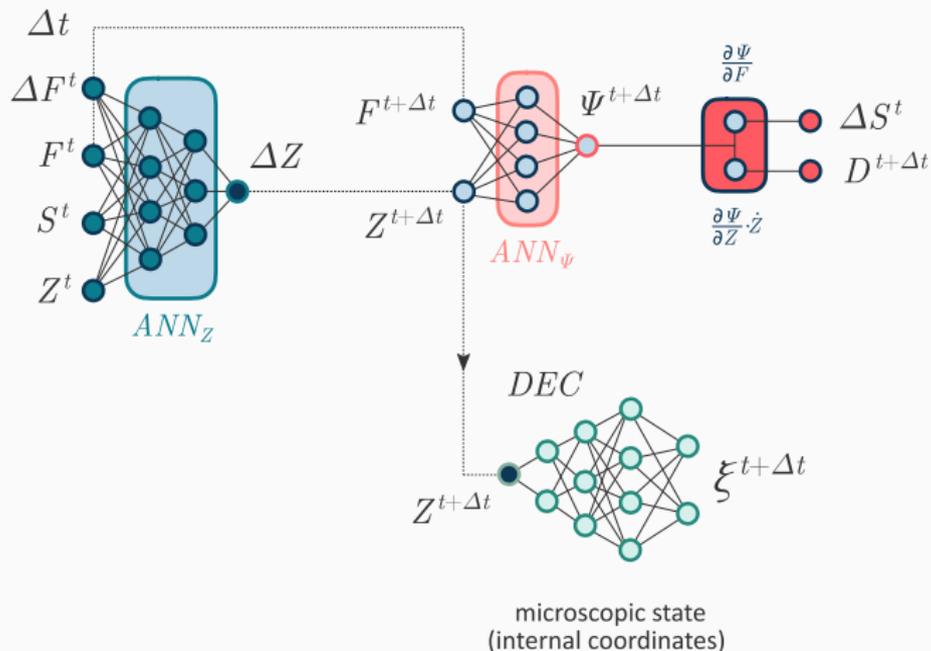
THERMODYNAMICS-BASED ARTIFICIAL NEURAL NETWORKS



$$\ell = \lambda^\xi [\ell(\xi_*^t) + \ell(\xi_*^{t+\Delta t})] + \lambda^{\nabla_{ED}^t} \ell(\nabla_{ED}^t) + \lambda^\Psi \ell(\Psi) + \lambda^{\Delta S} \ell(\Delta S) + \lambda^D \ell(D)$$

(Masi et al. arXiv preprint arXiv:2108.13137, 2021)

THERMODYNAMICS-BASED ARTIFICIAL NEURAL NETWORKS



$$\ell = \lambda^{\dot{Z}} \ell(\dot{Z}) + \lambda^\Psi \ell(\Psi) + \lambda^{\Delta S} \ell(\Delta S) + \lambda^D \ell(D)$$

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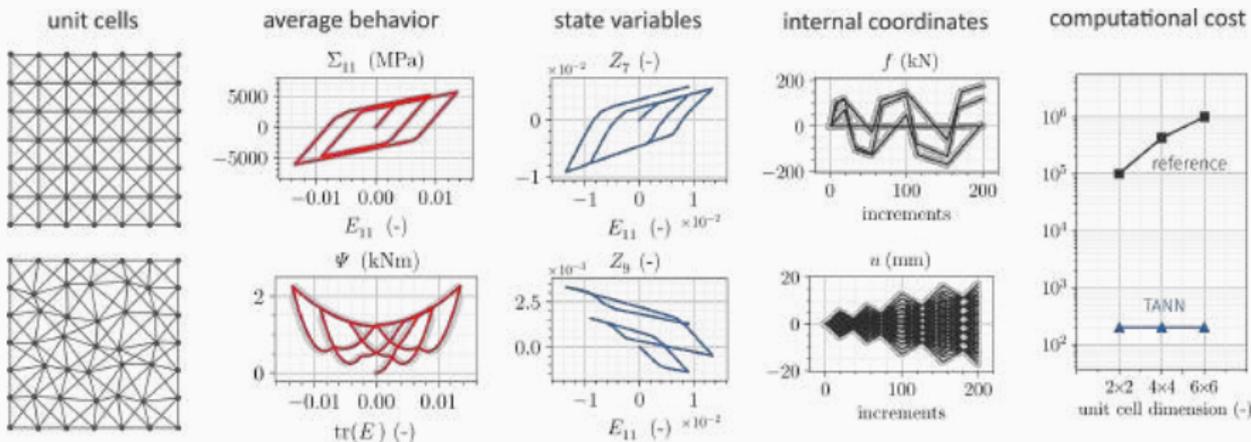
BENCHMARKS FOR LATTICE STRUCTURES

Inelastic lattice cells

$$\Sigma = \langle \sigma \rangle = \frac{1}{|V|} \sum_k^{N_{nodes}^p} t_k \otimes y_{\Delta k}^p$$

$$E = \langle \varepsilon \rangle = \frac{1}{|V|} \int_{\partial V} u \otimes x ds$$

- training on micromechanical datasets generated from random loading paths



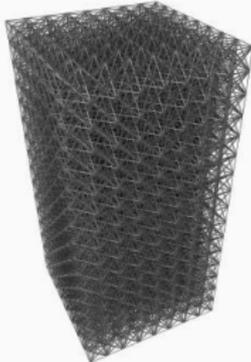
ACCURACY $\geq 99\%$

—

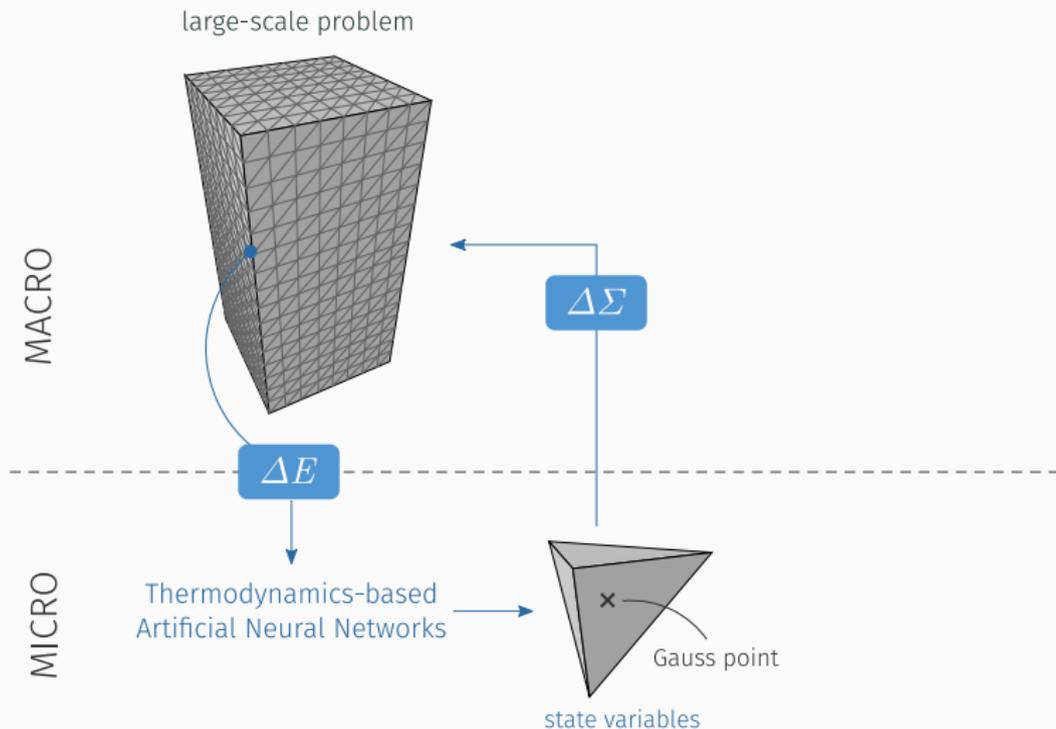
ACCELERATION $\approx 10^3$

THE FEM \times TANN APPROACH

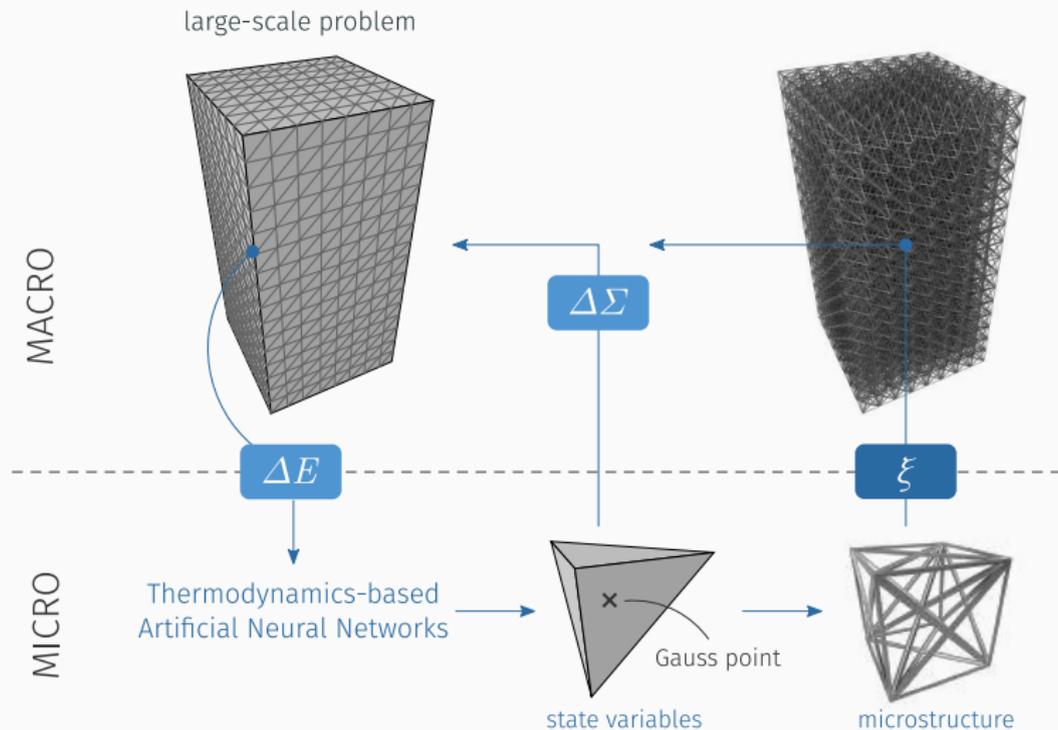
large-scale problem



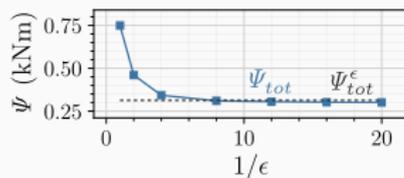
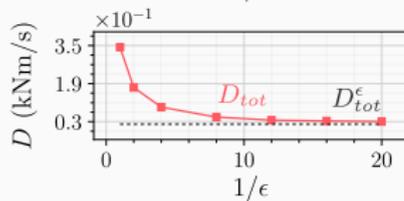
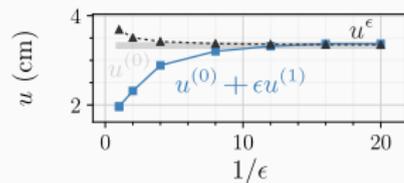
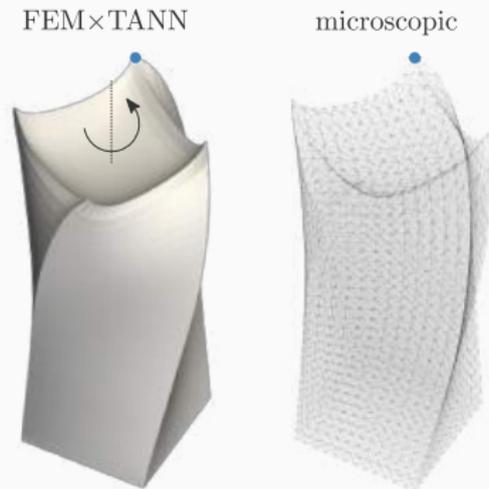
THE FEM×TANN APPROACH



THE FEM×TANN APPROACH



APPLICATION TO LARGE-SCALE PROBLEMS



APPLICATION TO LARGE-SCALE PROBLEMS



$\Omega = 0^\circ$



$\Omega = 30^\circ$



$\Omega = 0^\circ$



$\Omega = 30^\circ$



$\Omega = 50^\circ$



$\Omega = 30^\circ$



$\Omega = 0^\circ$



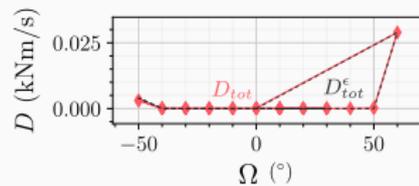
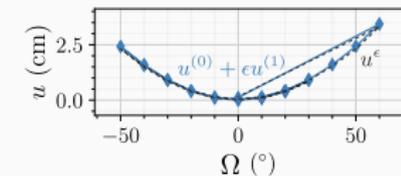
$\Omega = 30^\circ$



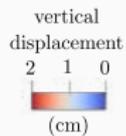
$\Omega = 50^\circ$



$\Omega = 60^\circ$



unloading



APPLICATION TO LARGE-SCALE PROBLEMS



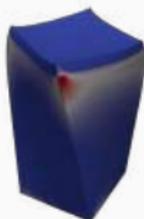
$\Omega = 0^\circ$



$\Omega = 30^\circ$



$\Omega = 0^\circ$



$\Omega = 30^\circ$



$\Omega = 50^\circ$



$\Omega = 30^\circ$



$\Omega = 0^\circ$



$\Omega = 30^\circ$



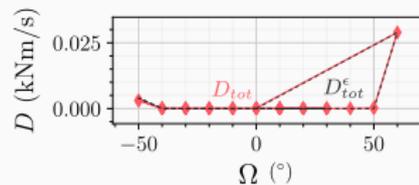
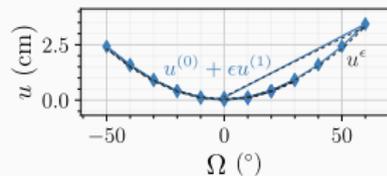
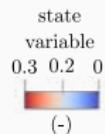
$\Omega = 50^\circ$



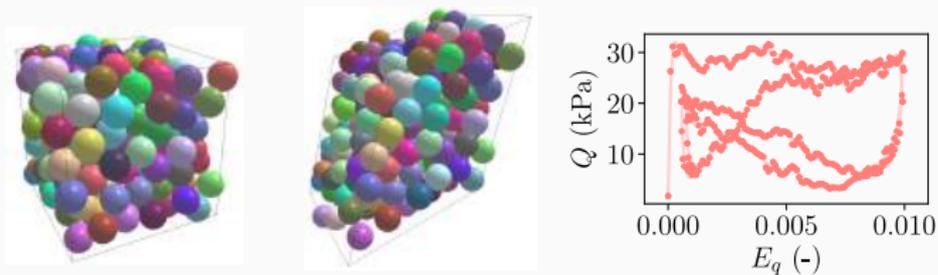
$\Omega = 60^\circ$



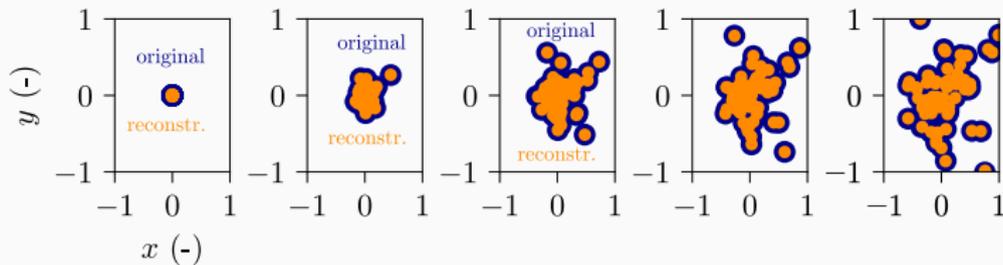
unloading



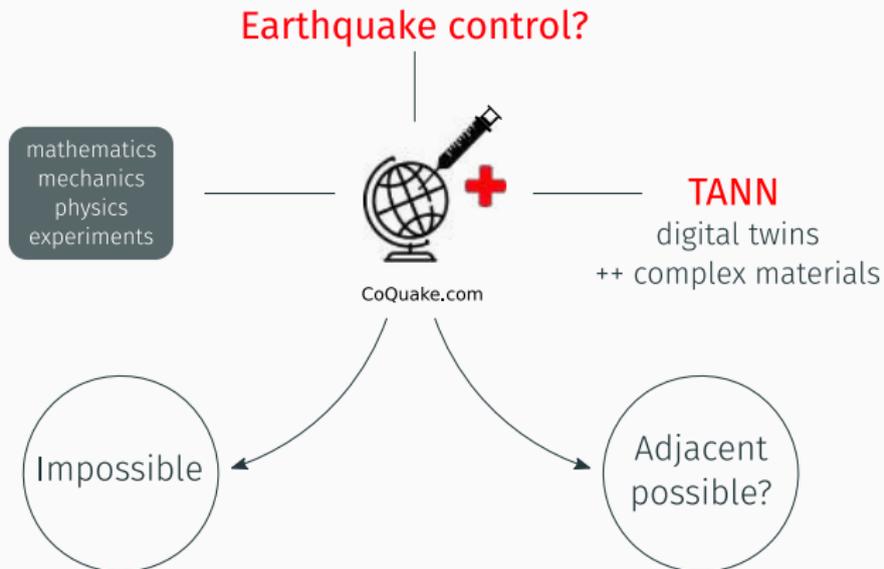
APPLICATION TO GRANULAR MEDIA



encoding the microscopic state space



CONCLUSIONS



ACKNOWLEDGMENTS

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Ioannis Stefanou



Alexandros Stathas



Farah Rabie

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- Masi et al (2021) Thermodynamics-based Artificial Neural Networks for constitutive modeling, *J Mech Phys Solids*. doi: [10.1016/j.jmps.2020.104277](https://doi.org/10.1016/j.jmps.2020.104277)

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THANK YOU!



THE UNIVERSITY OF
SYDNEY



Can we tame earthquakes?

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